Information Storage Capacity of Interacting Particle Systems

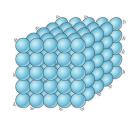
Ziv Goldfeld

Cornell University

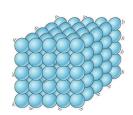
Collaborators: Guy Bresler and Yury Polyanskiy

Beyond IID in Information Theory 8

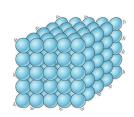
Nov. 13th, 2020



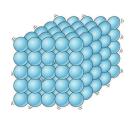
Writing data



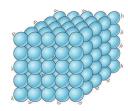
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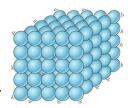
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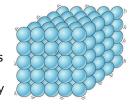
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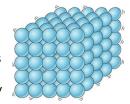


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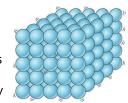
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Distilling notion of storage from particular technology

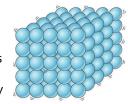
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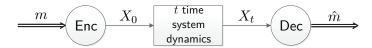
- Distilling notion of storage from particular technology
- Capturing interparticle interaction and system's dynamics

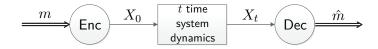
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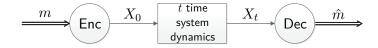


Goal: Study information storage capacity while:

- Distilling notion of storage from particular technology
- Capturing interparticle interaction and system's dynamics
- How much data can be stored and for how long?

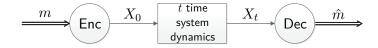






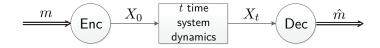
Stochastic Ising Model:

• Graph $(\mathcal{V}, \mathcal{E})$: topology of the storage medium.



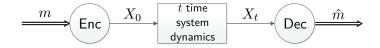
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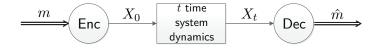


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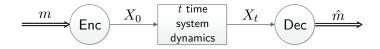
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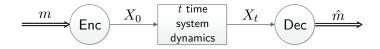
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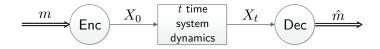
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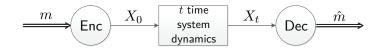
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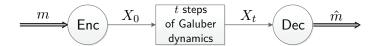
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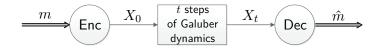
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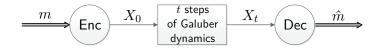
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Cold $(\beta \text{ large}) \implies \text{Strong interactions}$

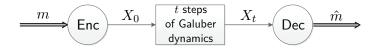




Information Capacity:

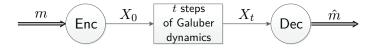


Information Capacity:
$$I_n^{(\beta)}(t) := \max_{P_{X_0}} I(X_0; X_t)$$



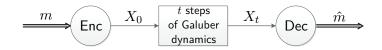
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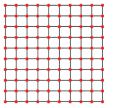
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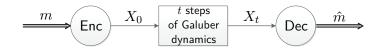
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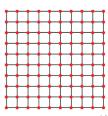
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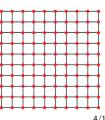
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$$\xrightarrow{\hspace*{1cm}} \text{Enc} \xrightarrow{\hspace*{1cm}} X_0 \xrightarrow{\hspace*{1cm}} \text{of Galuber } \\ \text{dynamics} \xrightarrow{\hspace*{1cm}} X_t \xrightarrow{\hspace*{1cm}} \text{Dec} \xrightarrow{\hspace*{1cm}} \hat{m} \xrightarrow{\hspace*{1cm}}$$

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- **Warm:** n-fold DM BSC $\left(\frac{1}{2} + o(1)\right)$ after t = O(n).
- **Cold:** Can interactions (memory) help?



Majority Update:

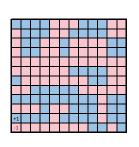
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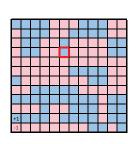
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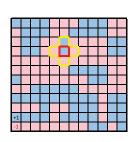
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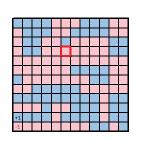
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Zero-Temperature Dynamics $(\beta \to \infty)$

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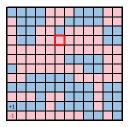
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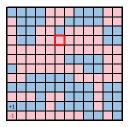


Operation Domain coarsening: Monochrom. clusters shrink/grow/split/coalesce

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Q1: What (if anything) can be stored for infinite time?

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For the zero-temp. SIM on $\sqrt{n} \times \sqrt{n}$ grid $I_n(\infty) := \lim_{t \to \infty} I_n(t) = \Theta(\sqrt{n})$

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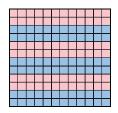
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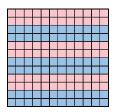
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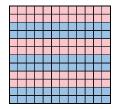
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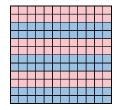
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For the zero-temp. SIM on $\sqrt{n} \times \sqrt{n}$ grid $I_n(\infty) := \lim_{t \to \infty} I_n(t) = \Theta(\sqrt{n})$

Achievability

- Stable Configurations: $\sigma \in \Omega$ is *stable* if $P(\sigma, \sigma) = 1$ (ground states).
- All 2-striped config. are stable.
- $\text{ \# Stripes} = 2^{\Theta(\sqrt{n})} \text{ \& } X_0 \sim \mathsf{Unif}\big(\{\mathsf{Stripes}\}\big) \Longrightarrow \ I_n(\infty) = \Omega(\sqrt{n})$

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$$\Longrightarrow \lim_{t \to \infty} \mathbb{P}(X_t \in \{\mathsf{Stripes}\}) = 1$$

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Converse:

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$$\Longrightarrow \lim_{t o\infty} \mathbb{P}ig(X_t \in \{\mathsf{Stripes}\}ig) = 1 \implies oldsymbol{I_n(\infty)} = oldsymbol{O}(\sqrt{n})$$

Q2: Can we do better than \sqrt{n} for finite superlinear t?

Theorem (G.-Bresler-Polyanskiy'19)

Let a(n) = o(n). Then $\exists c > 0$ s.t. $I_n(t) = \Omega\left(\frac{n}{a(n)}\right)$, $\forall t \leq c \cdot a(n) \cdot n$.

Theorem (G.-Bresler-Polyanskiy'19)

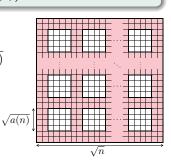
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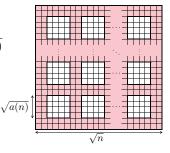
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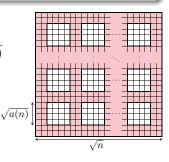
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- Separate by all-minus 2-strips



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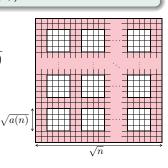
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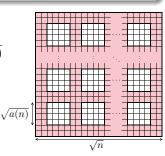


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• Continuous-Time: Updates according to i.i.d. Poiss(1/n) clocks.

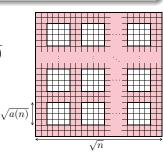
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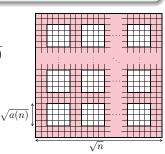
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⇒ Non-interacting portions are independent.

Tensorization: $I_n^{(c)}(t) \geq K \cdot \max_{p_1} I\left(\left[X_0^{(c)}\right]_1; \left[X_t^{(c)}\right]_1\right)$

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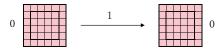
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Subsquare:



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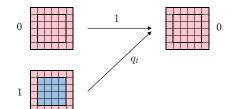


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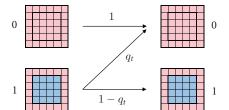


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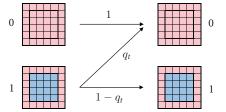
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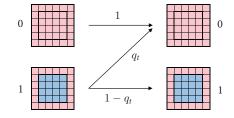


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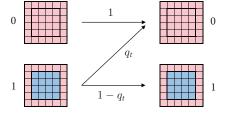


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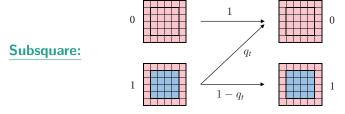
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$$\Longrightarrow I_n^{(c)}(c \cdot a(n) \cdot n) \ge C \cdot K = \Omega\left(\frac{n}{a(n)}\right)$$

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We've Seen:

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Further Questions:

① Upper bounds better than n for $t < \infty$?

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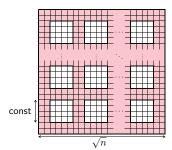
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- Improved scheme for superlinear time?
 - ▶ Nesting infinitely many sub-squares with vanishing growth rates.

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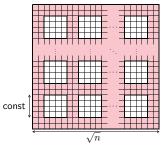


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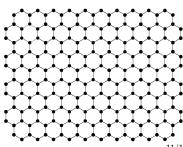
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Honeycomb Lattice (no external field):



Grid with External Field:

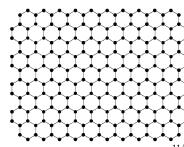
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Honeycomb Lattice (no external field):

• $deg(v) = 3, \forall v \text{ in interior}$



Grid with External Field:

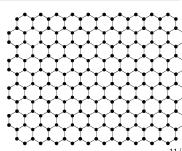
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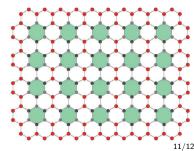
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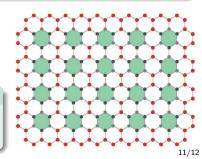
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12/12

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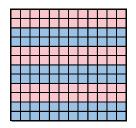
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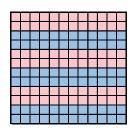
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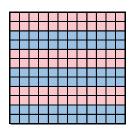
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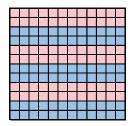
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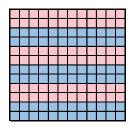
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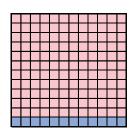
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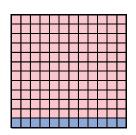
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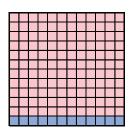
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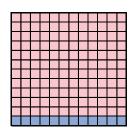
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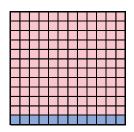


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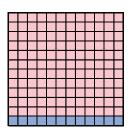
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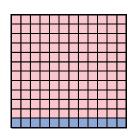
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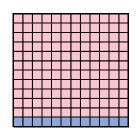


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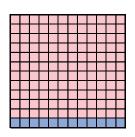
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Interleaved Dynamics: 2 types of flips

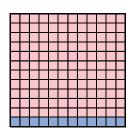


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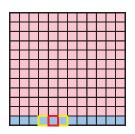


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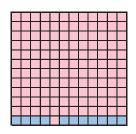


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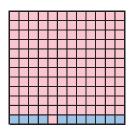
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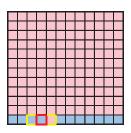
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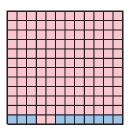
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• Initially chain stays close to X_0 w/ occasional sprinkles

ℜ Pluses may spread out above bottom stripe

Fix: Prohibit minus-spins from flipping (speedup)

- **® Interleaved Dynamics:** 2 types of flips
 - ► **Sprinkle:** Flip w/ all-plus horizontal neighbors
 - Erosion: Flip w/ at least one minus horizontal neighbor

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Expected Behavior:

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- After sufficiently many sprinkle, drift driven by erosion
- \implies Dominate $\{X_t\}_t$ by a phase-separated dynamics

• Consider continuous-time dynamics (i.i.d. Poisson clocks at each site)

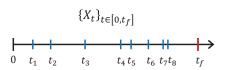
- Consider continuous-time dynamics (i.i.d. Poisson clocks at each site)
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- \bullet Define new dynamics $\{\tilde{X}_t\}_{t\in[0,2t_f]}$ with first 2k clock rings and flips

$$\tau_{j} = \begin{cases} t_{j}, & j \in [k] \\ t_{j-k} + t_{f}, & j \in [k+1:2k] \end{cases}, \ u_{j} = \begin{cases} v_{j}, & j \in [k] \\ v_{j-k}, & j \in [k+1:2k] \end{cases}$$

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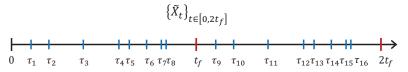
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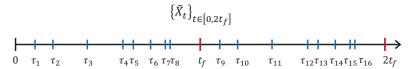
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$$\begin{cases} X_{t} \}_{t \in [0,t_{f}]} \\ 0 & t_{1} \quad t_{2} \quad t_{3} \quad t_{4} t_{5} \quad t_{6} t_{7} t_{8} \quad t_{f} \end{cases}$$

$$0 \quad \tau_{1} \quad \tau_{2} \quad \tau_{3} \quad \tau_{4} \tau_{5} \quad \tau_{6} \tau_{7} \tau_{8} \quad t_{f} \quad \tau_{9} \quad \tau_{10} \quad \tau_{11} \quad \tau_{12} \tau_{13} \quad \tau_{14} \tau_{15} \tau_{16} \quad 2t_{f} \end{cases}$$

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Blocking Rule:

$$\begin{split} & \big\{ \tilde{X}_t \big\}_{t \in \left[0, 2t_f\right]} \\ & \\ \downarrow & \\ 0 \quad \tau_1 \quad \tau_2 \qquad \tau_3 \qquad \tau_4 \tau_5 \quad \tau_6 \quad \tau_7 \tau_8 \qquad t_f \quad \tau_9 \quad \tau_{10} \qquad \tau_{11} \quad \tau_{12} \tau_{13} \, \tau_{14} \tau_{15} \, \tau_{16} \quad 2t_f \end{split}$$

Blocking Rule:

① For $t < t_f$ allow only sprinkle flips (wrt original $\{X_t\}_{t \in [0,t_f]}$)

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Observations:

 \bullet Erosion flips in $\{X_t\}_{t\in[0,t_f]}$ \implies Erosion flips in $\{\tilde{X}_t\}_{t\in[t_f,2t_f]}$

$$\left\{ \tilde{X}_{t} \right\}_{t \in \left[0, 2t_{f}\right]}$$

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Observations:

- ullet Erosion flips in $\{X_t\}_{t\in[0,t_f]}$ \Longrightarrow Erosion flips in $\{\tilde{X}_t\}_{t\in[t_f,2t_f]}$
- Erosion flip rates in $\{\tilde{X}_t\}_{t\in[t_f,2t_f]}$ are faster.

$$\left\{ \widetilde{X}_{t} \right\}_{t \in \left[0, 2t_{f}\right]}$$

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- Erosion flip rates in $\{\tilde{X}_t\}_{t\in[t_f,2t_f]}$ are faster.
- \Longrightarrow New dynamics is a speedup: $\mathbb{E}N^{(+)}(t_f) \geq \mathbb{E}\tilde{N}^{(+)}(2t_f)$

Sprinkle Analysis $[0,t_f]$: Ends w/ runs of '+'s separated by '-' sprinkles

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Q: What is the typical length of a run (contig) & how many of them?

• Approx. bottom stripe sites by i.i.d. $Exp(p_{\beta}), p_{\beta} \triangleq \mathbb{P}(Sprinkle)$

Sprinkle Analysis $[0,t_f]$: Ends w/ runs of '+'s separated by '-' sprinkles



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- ullet Approx. $\mathsf{L}_i=$ 'Length of Contig i' by $\mathsf{Geo}(p_\beta^{-1})$

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- Show $\mathbb{E}[\mathsf{Number} \ \mathsf{of} \ \mathsf{contigs} \ \mathsf{of} \ \mathsf{this} \ \mathsf{length}] \gtrsim \frac{\sqrt{n}}{2-p_\beta}$

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Erosion Analysis $(t_f, 2t_f]$: Contig eaten w/ speed $\phi_\beta \triangleq \frac{e^\beta}{e^\beta + e^{-\beta}}$ (2 sides)

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$$\implies$$
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• Show latter probability is small and conclude proof